



On generalized quasi-conformal like recurrent manifolds

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Abstract. The object of the present paper is to introduce a generalized class of recurrent manifold called quasi-conformal like recurrent manifolds. Some geometric properties of generalized quasi-conformal like recurrent manifolds are obtained. The application of such a manifold in general relativity are investigated.

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Introduction: It is well known that symmetric spaces play a great role in the differential geometry. The study of Riemann spaces was initiated by E. Cartan [8]. According to E. Cartan a Riemannian manifold is said to be locally symmetric if its Riemannian curvature tensor R of type $(0, 4)$ satisfies the condition $\nabla R = 0$ where ∇ denote the Levi-Civita connection. Later many authors have weakened this notion in different ways such as recurrent manifold by Walker [25], generalized recurrent by Dubey [12], quasi-generalized recurrent by Shaikh and Roy [21], hyper-generalized recurrent by Shaikh and Patra [20], semisymmetric manifold by Cartan [8], pseudosymmetric manifolds by Chaki [9], pseudosymmetric manifolds by Deszcz[11], weakly symmetric manifold by Tamássy and Binh [23]. Thereafter, a lot of research has been carried out in weakly symmetric manifold. For details, we refer to [1], [2], [3], [4], [5], [7] and the references there in.

In tune with Yano and Sawaki [24], recently Baishya and Chowdhury [6] introduce a new type of tensor field, named generalized quasi-conformal curvature tensor. The beauty of generalized quasi-conformal curvature tensor lies in the fact that it has the flavor of Riemannian curvature R , conformal curvature tensor C , conharmonic curvature tensor \bar{C} , concircular curvature tensor E , projective curvature tensor P and m -projective curvature tensor H as special cases. The generalized quasi-conformal curvature tensor is defined as

$$\tilde{W}(X, Y)Z = \frac{n-2}{n} [\{1 - b + (n - 1)a\} - \{1 + (n - 1)(a + b)c\}]C(X, Y)Z$$

$$+[1 - b(n - 1)a]E(X, Y)Z + (n - 1)(b - a)P(X, Y)Z + \frac{n-2}{n}(c - 1)\{1 + (n - 1)(a + b)\}\bar{C}(X, Y)Z \quad (1.1)$$

For all $X, Y, Z \in \chi(M^n)$, the set of all vector field of the manifold M , where a, b and c are real constants. The generalized quasi-conformal curvature tensor \tilde{W} is reduced to be (1) Riemannian curvature tensor R , if $a=b=c=0$, (2) conformal curvature tensor C if $a = b = -\frac{1}{n-2}$, $c=1$, (3) conharmonic curvature tensor \bar{C} , if $a = b = -\frac{1}{n-2}$, $c=0$, (4) concircular curvature tensor E , if $a=b=0$ and $c=1$, (5) projective curvature tensor P , if $a = -\frac{1}{n-1}$, $b=0$, $c=0$ and (6) M -projective curvature tensor H , if $a=b = -\frac{1}{2(n-1)}$ and $c=0$, is introduced by G. P. Pokhariyal and R. S. Mishra [17], which is defined as follows

$$H(X, Y)Z = R(X, Y)Z - \frac{1}{2(n - 1)}[S(Y, Z)X - S(X, Z)Y + g(Y, Z)QX - g(X, Z)QY],$$

Where Q be the symmetric endomorphism of the tangent bundle corresponding to the Ricci tensor S .

For our convenience in this paper we renamed generalized quasi-conformal curvature \tilde{W} as quasi-conformal like curvature tensor.

Generalizing the notion of recurrent manifold, in this paper we introduce the notion of generalized quasi-conformal like recurrent manifold. A Riemannian manifold $M, n > 2$, is said to be generalized quasi-conformal like recurrent manifold if its quasi-conformal like curvature tensor \tilde{W} satisfies the following relation:

$$(\nabla_X \tilde{W})(Y, U, V, W) = A(X)\tilde{W}(Y, U, V, W) + \alpha(X)G(Y, U, V, W). \quad (1.2)$$

$$\text{Where } G(Y, U, V, W) = [g(U, V)g(Y, W) - g(Y, V)g(U, W)]$$

And A, α are non-zero 1-forms which is defined as $A(X) = g(X, \theta)$, $\alpha = g(X, \phi)$. Such an n -dimensional manifold is denoted by $(GQCLK)_n$.

The beauty of such $(GQCLK)_n$ space is that for suitable choice of 1-forms and the real constants a, b , and c we get different spaces such as

- (i) Generalized recurrent manifolds [12] (for $a=b=c=0$),
- (ii) Generalized conformally recurrent manifolds [16] (for $a=b = -\frac{1}{n-2}$, $c=1$),
- (iii) Generalized conharmonically recurrent manifolds [20](for $a=b = -\frac{1}{n-2}$, $c=0$)
- (iv) Generalized concircularly recurrent manifolds[13](for $a=b=0$ and $c=1$)
- (v) Generalized M -projectively recurrent manifolds[15] (for $a=b = -\frac{1}{2(n-1)}$ and $c=0$)

- (vi) Recurrent manifolds[25] (for $a=b=c=\alpha = 0$)
- (vii) Symmetric manifolds[8] (for $a=b=c=A = \alpha = 0$).

We organized this paper as follows; section 2 is concerned with some basic geometric properties of $(GQCLK)_n$. In section 3 we find a sufficient condition for a $(GQCLK)_n$ to be a quasi-Einstein manifold. In the last section we studied some application of such manifold.

2. Some geometric properties of $(GQCLK)_n$.

In this section, we consider a Riemann manifold (M^n, g) $n > 2$ which is generalized quasi-conformal like recurrent manifold. Now, contracting (1.2) we have

$$(2.1) \quad (\bar{b} + \frac{n}{n-1} \bar{c}) (\nabla_X S)(Y, W) - \left(\frac{\bar{d}}{n-2} + \frac{\bar{b}}{n} + \frac{\bar{c}}{n-1} \right) dr(X)g(Y, W) = A(X) \left[\left(\bar{b} + \frac{n}{n-1} \bar{c} \right) S(Y, W) - \left(\frac{\bar{d}}{n-2} + \frac{\bar{b}}{n} + \frac{\bar{c}}{n-1} \right) rg(Y, W) \right] + (n-1)\alpha(X)g(Y, W)$$

Where $\bar{a} = \frac{n-2}{2} [\{1 - b + (n-1)a\} - \{1 + (n-1)(a+b)\}c]$, $\bar{b} = [1 - b + (n-1)a]$,

$\bar{c} = (n-1)(b-a)$ and $\bar{d} = \frac{n-2}{2} (c-1)\{1 + (n-1)(a+b)\}$.

On simplification we get

$$(2.2) \quad (\bar{b} + \frac{n}{n-1} \bar{c}) (\nabla_X S)(Y, W) = [A(X)] \left(\bar{b} + \frac{n}{n-1} \bar{c} \right) S(Y, W) + \left[(n-1)\alpha(X) - (rA(X) - dr(X)) \left(\frac{\bar{d}}{n-2} + \frac{\bar{b}}{n} + \frac{\bar{c}}{n-1} \right) \right] g(Y, W).$$

This leads the following:

Theorem 2.1. A $(GQCLK)_n$ is a generalized Ricci recurrent manifold [14] provided $(n-1)\bar{b} + n\bar{c} \neq 0$.

Taking contraction over Y and W in (2.1) we obtain

$$(2.3) \quad -\frac{n}{n-2} \bar{d} dr(X) = -\frac{n}{n-2} \bar{d} r A(X) + (n-1)n\alpha(X), \text{ i.e.,}$$

$$(2.4) \quad dr(X) = r A(X) - \frac{(n-1)(n-2)}{\bar{d}} \alpha(X).$$

If we take the scalar curvature of $(GQCLK)_n$ space is non-zero constant, then we have $dr(X)=0$ and from (2.4) we have

$$(2.5) \quad r A(X) - \frac{(n-1)(n-2)}{\bar{d}} \alpha(X) = 0$$

This gives the following:

Theorem 2.2. If the scalar curvature of a $(GQCLK)_n$ is non-zero constant, then the 1-forms are related by the expression (2.5).

Theorem 2.3. If the scalar curvature of a $(GQCLK)_n$ is non-zero constant, then the 1-forms are co directional provided $\bar{d} \neq 0$.

Again contraction of (2.1) over X and W yields

$$(2.6) \quad \left[\frac{n-2}{2} \left(\bar{b} + \frac{n}{n-1} \bar{c} \right) - \frac{\bar{d}}{n-2} \right] dr(Y) = \left(\bar{b} + \frac{n}{n-1} \bar{c} \right) A(QY) - \left(\frac{\bar{d}}{n-2} + \frac{\bar{b}}{n} + \frac{\bar{c}}{n-1} \right) rA(Y) + (n-1)\alpha(Y),$$

where $A(QX) = g(QX, \theta)$.

Using (2.4) in (2.6) we have

$$(2.7) \quad r = \frac{2A(QX)}{A(X)} + \frac{2(n-1)(n-2)^2 \alpha(X)}{A(X)n \bar{d}} \quad \text{provided } (n-1)\bar{b} + n\bar{c} \neq 0.$$

Thus we can state the following:

Theorem 2.4. In a $(GQCLK)_n$ the scalar curvature exist provided $(n-1)\bar{b} + n\bar{c} \neq 0$, $\bar{d} \neq 0$ and is given by (2.7).

Again if the scalar curvature $r = 0$ then we have

$$(2.8) \quad A(QY) = -\frac{(n-1)(n-2)^2 \alpha(Y)}{A(X)n \bar{d}}$$

Theorem 2.5. If the scalar curvature of a $(GQCLK)_n$ is vanishes, then the 1-forms are related by the expression (2.8).

If a $(GQCLK)_n$ is a Ricci-symmetric, then we have $\nabla S=0$, $dr = 0$ and r is constant. So from (2.2) we get

$$(2.9) \quad A(X) \left(\bar{b} + \frac{n}{n-1} \bar{c} \right) S(Y, W) + \left[(n-1)\alpha(X) - rA(X) \left(\frac{\bar{d}}{n-2} + \frac{\bar{b}}{n} + \frac{\bar{c}}{n-1} \right) \right] g(Y, W) = 0.$$

On simplification we have

$$(2.10) \quad S(Y, W) = -\frac{\left[(n-1)\alpha(X) - rA(X) \left(\frac{\bar{d}}{n-2} + \frac{\bar{b}}{n} + \frac{\bar{c}}{n-1} \right) \right] g(Y, W)}{A(X) \left(\bar{b} + \frac{n}{n-1} \bar{c} \right)} \quad \text{provided } (n-1)\bar{b} + n\bar{c} \neq 0.$$

This can be written as

$$S(Y, W) = \mu g(Y, W), \quad \text{where } \mu = -\frac{\left[(n-1)\alpha(X) - rA(X) \left(\frac{\bar{d}}{n-2} + \frac{\bar{b}}{n} + \frac{\bar{c}}{n-1} \right) \right]}{A(X) \left(\bar{b} + \frac{n}{n-1} \bar{c} \right)}, \text{ a}$$

scalar.

Thus we can state the following:

Theorem 2.6. A Ricci-symmetric $(GQCLK)_n$ is a Einstein manifold provided $(n-1)\bar{b} + n\bar{c} \neq 0$.

3. Sufficient condition for a generalized quasi-conformal like recurrent manifold to be a quasi Einstein manifold

In a $(GQCLK)_n$ the vector field θ defined by $g(X, \theta) = A(X)$ for any vector field X is said to be a concircular vector field [19] if the following equation is satisfied

$$(3.1) \quad (\nabla_X A)(Y) = \beta g(X, Y) + \varphi(X) A(Y), \quad \text{where } \beta \text{ is a non-zero scalar and } \varphi \text{ is a closed 1-form. If } \theta \text{ is a unit one, then (3.1) becomes}$$

$$(3.2) \quad (\nabla_X A)(Y) = \beta [g(X, Y) - A(X) A(Y)].$$

We consider that $(GQCLK)_n$ admits a unit concircular vector field defined by (3.1), where β is a non-zero constant. Proceeding in the same way in [15] in Riemann case we get the flowing result:

$$(3.3) \quad S(X, \theta) = (n - 1)\beta^2 A(X).$$

$$(3.4) \quad (\nabla_X S)(Y, \theta) = (n - 1)\beta^2 (\nabla_X A)(Y) - S(Y, \nabla_X \theta).$$

$$(3.5) \quad S(Y, \nabla_X \theta) = \beta[S(X, Y) - A(X)S(Y, \theta)].$$

Using (3.2) in (3.4) we have

$$(3.6) \quad (\nabla_X S)(Y, \theta) = (n - 1)\beta^3 [g(X, Y) - A(X)A(Y)] - S(Y, \nabla_X \theta).$$

By virtue of (3.5) it follows from (3.6) that

$$(3.7) \quad (\nabla_X S)(Y, \theta) = (n - 1)\beta^3 [g(X, Y) - A(X)A(Y)] - \beta[S(X, Y) - A(X)S(Y, \theta)].$$

Using (3.3) in (3.7) we get

$$(3.8) \quad (\nabla_X S)(Y, \theta) = (n - 1)\beta^3 g(X, Y) - \beta S(X, Y).$$

From (2.2) we have

$$(3.9) \quad \left(\bar{b} + \frac{n}{n-1} \bar{c}\right) (\nabla_X S)(Y, \theta) = [A(X)] \left(\bar{b} + \frac{n}{n-1} \bar{c}\right) S(Y, \theta) + \left[(n - 1)\alpha(X) - (rA(X) - dr(X)) \left(\frac{\bar{d}}{n-2} + \frac{\bar{b}}{n} + \frac{\bar{c}}{n-1}\right)\right] g(Y, \theta).$$

Now using (3.8) and (3.3) in (3.9) we get

$$(3.10) \quad \left(\bar{b} + \frac{n}{n-1} \bar{c}\right) [(n - 1)\beta^3 g(X, Y) - \beta S(X, Y)] = A(X) \left(\bar{b} + \frac{n}{n-1} \bar{c}\right) (n - 1)\beta^2 A(Y) + \left[(n - 1)\alpha(X) - (rA(X) - dr(X)) \left(\frac{\bar{d}}{n-2} + \frac{\bar{b}}{n} + \frac{\bar{c}}{n-1}\right)\right] A(Y).$$

Now if the scalar curvature is constant, then $dr = 0$ and using (2.5) in (3.10)

$$(3.11) \quad S(X, Y) = \beta^2 (n - 1)g(X, Y) + \frac{r - (n-1)\beta^2}{\beta} A(X)A(Y) \quad \text{provided } (n-1)\bar{b} + n\bar{c} \neq 0.$$

Since β is non-zero constant (3.11) can be written as

$$S(X, Y) = l g(X, Y) + m A(X)A(Y) \quad \text{provided } (n-1)\bar{b} + n\bar{c} \neq 0.$$

Where $l = \beta^2(n-1)$ and $m = \frac{r - (n-1)\beta^2}{\beta}$ are two non-zero constant as β is non-zero constant. Hence the manifold is a quasi-Einstein manifold. This leads us the following:

Theorem 3.1. A $(GQCLK)_n$ with constant scalar curvature becomes a quasi Einstein manifolds if the associated vector field θ defined by $g(X, \theta) = A(X)$ is a unit one and if it is a unit concircular vector field whose associated scalar is a non-zero constant.

4. Applications of perfect fluid Ricci symmetric $(GQCLK)_n$ space time

This section deals with $(GQCLK)_n$ spacetime. A spacetime is a connected 4-dimensional Lorentzian manifold. Hence a generalized quasi conformal like recurrent spacetime is a 4-dimensional connected Lorentzian generalized quasi conformal like recurrent manifold in which the vector field θ associated to the 1-form A is a unit timelike vector field. Such spacetimes is called $(GQCLK)_n$ spacetimes. Here we consider a special type of spacetime which is called Ricci symmetric generalized quasi conformal like recurrent spacetime. A semi-Riemannian four-dimensional Ricci symmetric generalized quasi-conformal like recurrent manifold may similarly be defined by taking a Lorentz metric g with signature $(+,+,+,-)$. In this case we consider a Ricci symmetric $(GQCLK)_n$ spacetime with the timelike velocity vector field $g(\theta, \theta) = -1$. So, Theorem 2.6. will also hold in such a spacetime. The Einstein’s equation without cosmological constant is

$$(4.1) \quad S(X, Y) - \frac{r}{2}g(X, Y) = k T(X, Y).$$

Where k is the gravitational constant, T is the energy- momentum tensor of type(0,2). We consider a perfect fluid spacetime. Then the energy- momentum tensor is of the form

$$(4.2) \quad T(X, Y) = (\sigma + \rho)A(X)A(Y) + \rho g(X, Y) ,$$

where σ and ρ as the energy density and isotropic pressure of the fluid respectively.

By virtue of (4.2), (4.1) can be written as

$$(4.3) \quad S(X, Y) - \frac{r}{2}g(X, Y) = k [(\sigma + \rho)A(X)A(Y) + \rho g(X, Y)]$$

Taking contraction (4.3) over X and Y we get

$$(4.4) \quad r = k(\sigma - 3\rho).$$

Now setting $W = X = \theta$ and using (2.5) in (2.10) for $n=4$ we get

$$(4.5) \quad S(Y, \theta) = \frac{3 \alpha(\theta) [1 - \frac{2}{d}(\frac{\bar{d}}{2} + \frac{\bar{b}}{4} + \frac{\bar{c}}{3})] A(Y)}{(\bar{b} + \frac{4}{3}\bar{c})}$$

Again substituting $Y = \theta$ in (4.3) we get

$$(4.6) \quad S(X, \theta) = \frac{r}{2} A(X) - k \sigma A(X).$$

By virtue of (4.4) and (4.5) , (4.6) takes the form

$$(4.7) \quad \alpha(\theta) = - \frac{k(\sigma+3\rho)(\bar{b}+\frac{4}{3}\bar{c})}{6[1-\frac{2}{d}(\frac{\bar{d}}{2} + \frac{\bar{b}}{4} + \frac{\bar{c}}{3})]}$$

Again applying (4.7) in (4.5) we get

$$(4.8) \quad S(Y, \theta) = - \frac{k(\sigma+3\rho)}{2} A(Y)$$

Putting $Y = \theta$ in (4.8) we have

$$(4.9) \quad S(\theta, \theta) = \frac{k(\sigma+3\rho)}{2}$$

From (4.3) and (4.4) we get

$$(4.10) \quad S(X, Y) = k \left[(\sigma + \rho)A(X)A(Y) + \frac{(\sigma - \rho)}{2} g(X, Y) \right]$$

And hence

$$(4.11) \quad S(QX, Y) = k \left[(\sigma + \rho)A(QX)A(Y) + \frac{(\sigma - \rho)}{2} S(X, Y) \right]$$

Taking contraction (4.11) over X and Y we get

$$(4.12) \quad \|Q\|^2 = k \left[(\sigma + \rho)S(\theta, \theta) + \frac{(\sigma - \rho)}{2} r \right]$$

Using (4.4) and (4.9) we obtain

$$(4.13) \quad \|Q\|^2 = k^2(\sigma^2 + 3\rho^2).$$

If we take the square of the length of the Ricci operator of the perfect fluid $(GQCLK)_n$ spacetime be $\frac{1}{3} r^2$, where r is the scalar curvature of the spacetime, then from (4.13), we have

$$(4.14) \quad \frac{1}{3} r^2 = k^2(\sigma^2 + 3\rho^2).$$

Which yields by virtue of (4.4) that

$$(4.15) \quad k^2(\sigma + 3\rho)\sigma = 0.$$

Since $\sigma + 3\rho \neq 0$ and $k \neq 0$, it follows from (4.15) that $\sigma = 0$, which is not possible as when the pure matter exists, σ is always greater than zero. Hence the spacetime under consideration can not contain pure matter. Thus we can state the following:

Theorem 4.1. A perfect fluid Ricci symmetric $(GQCLK)_n$ spacetime satisfying Einstein's field equation without cosmological constant can not contain pure matter if the square of the length of the Ricci operator is $\frac{1}{3} r^2$.

We know that if the Ricci tensor S of type (0,2) of the spacetime satisfies the condition [18]

$$(4.16) \quad S(X, X) > 0,$$

For every timelike vector field X , then (4.16) is called timelike convergence condition. Now we determine the sign of the pressure in such a spacetime without pure matter. Hence for $\sigma = 0$, (4.4) yields

$$(4.17) \quad r = -3\rho k.$$

Hence from (4.2) we get

$$(4.18) \quad T(\theta, \theta) = \sigma = 0.$$

Thus from (4.1) and (4.17) it follows that

$$\rho = \frac{2}{3k} S(\theta, \theta)$$

Since $S(\theta, \theta) > 0$, it follows from above relation that $\rho > 0$. Thus we can state the following:

Theorem 4.2. A perfect fluid Ricci symmetric $(GQCLK)_n$ spacetime having no pure matter, satisfying Einstein's field equation without cosmological constant and if its Ricci tensor follows the timelike convergence condition, the pressure of the fluid is positive.

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